

Eigenclustering-Length Exchange Theorems.

III. The Bottom Series: Straight- P_n

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Abstract

We have found a set of N -body problem Theorems, or alternatively for N -simplex configurations. Each is for a weighted sum of squares of those eigenclustering (= relative Jacobi) magnitudes which are not already (Lagrange) relative separations in terms of those which are. By which these are Eigenclustering Length-Exchange Theorems (ELETs). In the current Series, we restrict ourselves to equal-mass such, with reference to placing masses at the vertices; this is the most Geometrically-natural case. So we are more specifically considering EMELETs.

Eigenclustering networks become ambiguous for $N \geq 4$. In fact, these are unlabelled rooted binary tree (URBT)-valued. Though for the purpose of systematically naming our Theorems, it is more convenient to defoliate these once. Thus passing to the unlabelled rooted AMB trees: at-most binary.

Then the minimum nontrivial $N = 3$ EMELET returns Apollonius' Theorem for the median length in terms of sides data as the 2-path P_2 case. While the P_3 -bent- (alias 4-body H-)EMELET returns Euler's 4-Body Theorem for the Newton length in terms of separations data. In the current article, we show that trees smaller than P_2 do not contribute any ELETs. And, more substantially, we provide the explicit form of the EMELET for all straight P_n , $n \geq 2$. Where $n := N - 1$.

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1 Introduction

We recently gave new proofs [26, 29] of Apollonius' Theorem [1, 9, 18, 33] and Euler's 4-Body Theorem [3, 12, 13, 16, 17, 23] These use just centre of mass (CoM) and moment of inertia (MoI) concepts: first and second moments. By which they readily generalize to give corresponding Theorems for each eigenclustering network supported by each N -body problem in arbitrary dimension. For a weighted sum of squares of those eigenclustering magnitudes which are not already relative separations in terms of those which are.

Eigenclustering networks [32, 33] are more widely known as types of *Jacobi coordinate system* [8, 10, 11, 14]. In Sec 2, we set up separation coordinates and then the particular choice of eigenclustering coordinates that we require. Apollonius' Theorem in its original context is for a triangle: how to compute the length of a median from sides data. Via e.g. our proof, this generalizes to being a 3-body problem result, thus in particular holding in 1- d as well. The 3-body problem supports just the one eigenclustering network. While the most well-known subcase of Euler's 4-Body Theorem – Euler's Quadrilateral Theorem – is for the Newton length [2, 5, 6, 7, 17, 19] between the midpoints of the diagonals, in terms of separations data. Via e.g. our proof,¹ this generalizes to being a 4-body problem result, thus holding in 1- and 3- d as well. It corresponds to the H -eigenclustering network for the 4-body problem, wherein one alternative is supported: the K -eigenclustering network.

In this manner, our Theorems are Eigenclustering Length-Exchange Theorems (ELETs). More specifically, the current Series considers EMELETs. With reference to placing equal masses at the vertices, which is the most Geometrically-natural modelling situation.

Eigenclustering networks moreover correspond to [14, 32] the unlabelled rooted binary trees (*URBT*) [24]. We subsequently index eigenclustering networks, and thus our Theorems, by the corresponding trees. We

¹This is not the only way of obtaining this generalization, however; see e.g. [12].

explained this correspondence in Article I. As well as how it is more convenient to use the corresponding AMB trees – at most binary – that result from defoliating the URBT. See [32] for many further details about this.

From this point of view, the 3-body problem’s eigenclustering network corresponds to the 2-path tree P_2 . And the 4-body problem’s H versus K eigenclustering-network ambiguity corresponds to the bent versus straight P_3 rooting ambiguity. See Fig 1. These respectively support Apollonius’ Theorem as the 3-EMELET. Euler’s 4-Body Theorem as the H-EMELET. And the corresponding K-EMELET unveiled in [30].

We explain in Sec 3 how the space of nontrivial ELETs is a slight truncation of the URBT, with the Apollonius P_2 as its bottom element. We then present our new straight n -path Theorem in Sec 4. For $N \geq 5$, this corresponds to the $K(N)$ -eigenclustering generalization of $N = 4$ ’s K-eigenclustering. In Sec 5 we check that this recovers Apollonius’ Theorem (P_2 case) and my Jacobi-K Theorem (P_3 -straight case). And then finally give the explicit form for the next smallest case: the P_4 -straight eigenclustering for the 5-body problem. This eigenclustering is displayed in the last row of Fig 1.

2 Eigenclustering coordinates and their masses

2.1 Position and separation levels of structure

Structure 0 We denote the position vectors of our model’s points-or-particles by \mathbf{q}^I , $I = 1$ to N .

Structure 1 We denote the (*Lagrange*) (*relative*) *separation vectors* between them by \mathbf{r}^{IJ} , $J \geq I$. We also re-index these 2-index combinations by the following single index.

$$S = 1 \text{ to } \binom{N}{2} = \frac{N(N-1)}{2} = \dim(\text{separation space [27]}) . \quad (1)$$

We also use s^S to denote the corresponding magnitudes: the *separations* themselves. For a particular planar convex Geometrical figure’s realization, separations can be subdivided as follows. Into *sides* $a^I = s^I$, $I = 1$ to N . And *diagonals* d^D . With index

$$D = 1 \text{ to } \frac{N(N-1)}{2} - N = \frac{N(N-3)}{2} ; \quad (2)$$

see Fig 1.a) for the first few instances. Thereby, the first N S^S have ‘dual nationality’ as sides.

Structure 1’ With the \mathbf{r} not all being LI for $N \geq 3$, we introduce some relative space basis for them that we denote by $\tilde{\mathbf{r}}^i$. With

$$i = 1 \text{ to } n := N - 1 = \dim(\text{relative space [27]}) .$$

Remark 1 For $N \geq 3$, the inertia quadric is not however diagonal with respect to such a basis [27]. To acquire this property, we pass to the following.

2.2 Eigenclustering vectors, alias relative Jacobi vectors

Structure 1’’ *Eigenclustering vectors* [27], alias *relative Jacobi vectors* [4, 8, 11, 20, 33] are an alternative basis for the relative space information that is diagonalizing. We present this here in the case in which the underlying point-or-particle masses are equal.

Examples 0 to 2 For $N \leq 2$, the eigenclustering concept is unnecessary.

Example 3 For $N = 3$, they are uniquely specified up to 3 relabellings. Fixed by which side we choose to be the base of the triangle.



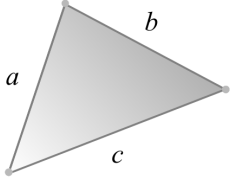
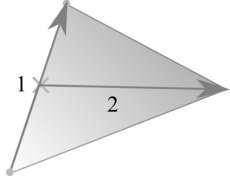
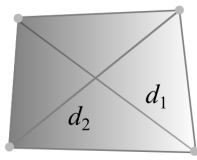
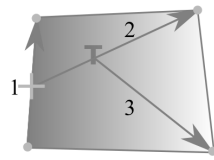
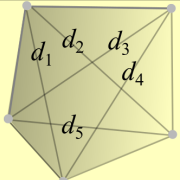
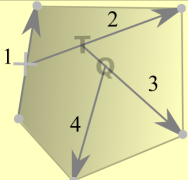
N	a) Separations	b) $K(N)$ -Eigencustering vectors	c) Corresponding AMB tree graphs
2			© 2024 Dr E. Anderson 1 • 1-path P_1
3			1 • 2 • 2-path P_2
4			1 • 2 • 3 • P_3 -straight
5			1 • 2 • 3 • 4 • P_4 -straight

Figure 1:

Example 4 For $N \geq 4$, there are further eigencustering network ambiguities. Starting with the H versus K ambiguity for the 4-body problem. Which corresponds to the bent versus straight P_3 rooting ambiguity. Indeed, the possible eigencustering networks for the N -body problem correspond to the unlabelled trees on N vertices. In the current Article, we consider the P_n -straight alias $K(N)$ the 4-body problem's P_3 alias K . Fig 1 illustrates this and the previous Subsection's main notions up to $N = 5$.

2.3 Straight-path eigencusterings

Remark 1 For our path choice, each N extends the preceding N 's eigenbasis by one eigenvector. From the total CoM so far (X , T , Q , P ... in Fig 1) to the extra point-or-particle. Thereby, one can immediately write down the general N th case (Fig 2.a).

Structure 2 Every eigencustering network contains ≥ 1 relative separation. For path eigencusterings, it is precisely 1. Eigencustering line segments which are not separations are some kind of [transversal]. This is a term introduced in [28] to cover both co-transversals such as medians and transversals such as the Newton line in the H case. See Fig 2.a) for explicit formula for the P_n -straight cases' eigencustering vectors.

Naming Remark 1 Instead of naming individual transversals as we did in [29, 30] for H and K respectively, we now just label them by their edges along P_n -straight. Edge 1 is the 2-body separation. Edge 2 extends this to a 3-body eigencustering basis, and so on. See Fig 1.b).

Path eigencustering for the N -body problem									
Eigencustering = relative Jacobi coordinates								© 2024 Dr E. Anderson	
\bar{R}_1	$:=$	\bar{q}^B	$-$	\bar{q}^A	$\}$	2-body	$\}$	3-body	$\}$
\bar{R}_2	$:=$	\bar{q}^C	$-$	$\frac{1}{2}(\bar{q}^A + \bar{q}^B)$	$\}$			4-body	$\}$
\bar{R}_3	$:=$	\bar{q}^D	$-$	$\frac{1}{3}(\bar{q}^A + \bar{q}^B + \bar{q}^C)$	$\}$			5-body	$\}$
\bar{R}_4	$:=$	\bar{q}^E	$-$	$\frac{1}{4}(\bar{q}^A + \bar{q}^B + \bar{q}^C + \bar{q}^D)$	$\}$			\dots	$\}$
\vdots									
\bar{R}_n	$:=$	\bar{q}^N	$-$	$\frac{1}{n} \sum_{i=1}^n \bar{q}^i$	$\}$			\dots	$\}$
								N -body	
Corresponding eigencustering masses									
$\frac{1}{\mu_1}$	$=$	$\frac{1}{1} + \frac{1}{1}$	$=$	2	\Rightarrow	$\mu_1 =$	$\frac{1}{2}$	$\}$	2-body
$\frac{1}{\mu_2}$	$=$	$\frac{1}{2} + \frac{1}{1}$	$=$	$\frac{3}{2}$	\Rightarrow	$\mu_2 =$	$\frac{2}{3}$	$\}$	3-body
$\frac{1}{\mu_3}$	$=$	$\frac{1}{3} + \frac{1}{1}$	$=$	$\frac{4}{3}$	\Rightarrow	$\mu_3 =$	$\frac{3}{4}$	$\}$	4-body
$\frac{1}{\mu_4}$	$=$	$\frac{1}{4} + \frac{1}{1}$	$=$	$\frac{5}{4}$	\Rightarrow	$\mu_4 =$	$\frac{4}{5}$	$\}$	5-body
\vdots						\vdots			\dots
$\frac{1}{\mu_n}$	$=$	$\frac{1}{n} + \frac{1}{1}$	$=$	$\frac{n+1}{n}$	$=$	$\frac{N}{n}$	\Rightarrow	$\mu_n =$	$\frac{n}{n+1} = \frac{n}{N}$
									N -body

Figure 2:

Notational Remark 1 For path eigencusterings, the corresponding edge-labelled [transversal] length variables are

$$t_2, t_3, \dots, t_n. \quad (3)$$

Thus the *path-eigencustering* [transversal] *subspace* of relative space has the following dimension.

$$n - 1 = N - 2.$$

Let us also use

$$T_2, \dots, T_n \quad (4)$$

to denote their squares.

2.4 Eigencustering masses and the ‘Greek world’

Structure 1 Eigencustering masses are a subcase of reduced masses. We compute these in Fig 2.b) for $N = 2$ to 5 . The general case for this is obvious as well. Observe that equal particle masses *does not* guarantee equal reduced masses, eigencustering masses included.

Notational Remark 2 Eigencustering masses lead to corresponding mass-scaled length variables [27]. Let us use Greek letter versions of symbols to denote mass-weighted counterparts. Chief among these in the current Article is ℓ , which, unadorned, stands for total moment of inertia. With indices, this stands for the corresponding partial moments. Though we use furthermore α corresponding to side a and $\tau_{\bar{\ell}}$ for those corresponding to [transversals] $t_{\bar{\ell}}$.

Notational Remark 3 Finally let us use \mathcal{O}_p to denote the sum of the p th powers of the objects \mathcal{O} . [21, 26, 25, 29, 30] benefited from many simplifications from sequential use of squared variables and power sums. This usefulness largely recurs in the current Article.

3 How our result trivializes for $N \leq 2$: the 3 smallest binary trees

Remark 1 The system of equations that our Theorem follows from consists of the following.

1) The MoI \mathcal{L} 's partial moments expansion in our eigenclustering basis,

$$\mathcal{L} = \sum_{i=1}^n \mathcal{L}_i =: \mathcal{L}_1 . \quad (5)$$

2) And the separations-democratic radius of gyration (RoG) formula [27]: an equable expansion in the separations. Which, in the equal-masses case used in the current Series, is

$$(\text{RoG})^2 := R = \frac{1}{N^2} \sum_{\substack{I, J = 1 \\ I < J}}^N \|\bar{\mathbf{r}}^{IJ}\|^2 . \quad (6)$$

By Sec 2, these expansions contain

$$n := N - 1 \quad (7)$$

and

$$\frac{N(N-1)}{2} \quad (8)$$

terms respectively. We are to solve our system for the [transversal] lengths. For our Theorem is a device from computing some (perhaps weighted) sum of such lengths squared in terms of purely separation-length data. This fits the bill: firstly, a truer name is Apollonius' *Median-Length* Theorem [26]. Secondly, Euler's 4-Body Theorem can be taken to compute the length of the Newton line segment [29].

Question 1 So what happens for the cases even smaller than Apollonius' Theorem?

Example 2 For the 2-body problem, there is just the 1 separation and no [transversals]. At the level of counting,

$$\frac{2(2-1)}{2} = 1 = 2 - 1 .$$

So both expansions become trivial, in the sense that there is only 1 piece in each. And our system of equations degenerates to 2 copies of the same identity equation,

$$\mathcal{L} = \alpha . \quad (9)$$

And this is fine, since there are also no [transversals] to solve for in this case!

The 2-body problem is thus too simple to have an Eigenclustering Length-Exchange Theorem. This case is indexed by the

$$\text{pt} = \text{D}_1 = \text{P}_1 \text{ AMB tree} .$$

Which are respectively, its single-point point-cloud notation, its totally-disconnected graph notation and its 1-path reconceptualiation.

Example 1 For the 1-body problem, there are no separations or [transversals]. Indeed now both counts return 0 . This means that both of our expansions contain no terms. Our system has thus again degenerated to 2 copies of the same equation, which is now furthermore a zero equation:

$$\mathcal{L} = 0 . \quad (10)$$

And this is fine, since there are no [transversals] at all to solve for in this case! So now the partial moments expansion is even more trivial: it contains *no* partial moments.

This case is indexed by the same AMB tree as above. This corresponds to the only place where the binary-to-AMB defoliation fails to give an isomorphism [32]. For, on the one hand, defoliating P_3 -bent returns

the point. But, on the other hand, defoliating the point also returns the point, because now the input tree already has no leaves at all! Thus if discussing $N = 1, 2$ we need AMB labelling to distinguish between the two. I.e. the P_3 -straight AMB tree versus the pt AMB tree.

Example 0 For the 0-body problem, there are no points-or-particles, let alone separations or [transversals]. Now neither a MoI nor a RoG exist. So there are no equations at all! And this is fine, since there are no [transversals] to solve for in this case either... This case is indexed by

$$U : \text{whichever of the unpoint, untree or unpath incarnations of } \emptyset . \quad (11)$$

Structure 1 Our Theorems are thus indexed in more detail by

$$\mathfrak{T}ree_{2*} = \{U, D, P\} . \quad (12)$$

For $\mathfrak{T}ree_{2*}$ the arena of unlabelled rooted binary trees. On which restriction it is isomorphic to

$$\mathfrak{T}ree_{\leq 2*} = \{U, D, P\} . \quad (13)$$

For $\mathfrak{T}ree_{\leq 2*}$ the arena of unlabelled rooted AMB trees.

While the current article involves just

$$\mathfrak{p}ath = \{U, D, P\} \cong \mathbb{N}_0 = \{0, 1, 2\} \cong \mathbb{N} . \quad (14)$$

For $\mathfrak{p}ath$ the arena of unlabelled paths. Thereby, Apollonius's Theorem – as the P_2 case – is, from an Order-Theoretic [15, 24] point of view, the *bottom element* of our slightly truncated arenas of path and tree graphs.

4 The P_n -straight Eigencenter [Transversal]-lengths² Theorem

Theorem 1 [Anderson 2018]

a) (Mass-weighted rational form)

$$N \sum_{\bar{t}=2}^n \tau_{\bar{t}} = 2 \Sigma_1 - N \alpha . \quad (15)$$

b) (Geometrical [transversal]-lengths subject form)

$$2 \sum_{\bar{t}=2}^n \frac{\bar{t}}{\bar{t} + 1} T_{\bar{t}} = \frac{2}{N} S_1 - A . \quad (16)$$

Proof Our First Principle (A) is the partial MoI expansion of in H -coordinates.

Our Second Principle (B) is the separations-democratic RoG formula.

Squared variables and sum variables sequentially save us symbols.

Latin \longleftrightarrow Greek ‘translations’ (inter-conversions) must always remember to (un)deploy eigenclustering mw factors.

Everything else is in Fig 3.

Remark 1 In the Geometrical variables, the eigenclustering masses render placing the general case in rational form unwieldy. This can however be done for each small example of interest ([30] and Sec 5).

Corollary 1 a)

$$N \sum_{\bar{t}=2}^n \tau_{\bar{t}} = (2 - N) \alpha + 2 \sum_{\bar{S}=2}^{\frac{N(N-1)}{2}} \Sigma_{\bar{S}} . \quad (17)$$

Proof of P_N eigenclustering shape N -body analogue of Euler's 4-Body Theorem	
Plain = Latin world	Mass-weighted = Greek world
$(B) \quad N^2 R = S_1$	$(A) \quad l = l_1$ $(B') \quad N l = 2 \Sigma_1$
	$N l_1 = 2 \Sigma_1$
	$N \left(\alpha + \sum_{\bar{i}=2}^n \tau_{\bar{i}} \right) = 2 \Sigma_1$
	$\sum_{\bar{i}=2}^n \tau_{\bar{i}} = 2 \Sigma_1 - N \alpha$
$2 N \sum_{\bar{i}=2}^n \frac{\bar{t}}{\bar{t} + 1} T_{\bar{i}} = 2 S_1 - N A$	
by cluster-labelling-choice and dimension-independent means! \square	

Figure 3:

b)

$$\sum_{\bar{i}=2}^n \frac{\bar{t}}{\bar{t} + 1} T_{\bar{i}} = \frac{2 - N}{2N} A + \frac{1}{N} \sum_{\bar{S}=2}^{\frac{N(N-1)}{2}} S_{\bar{S}} \quad (18)$$

Proof This follows from also splitting the second sum in Fig 2's second column. And then cancelling contributions from our two bases' 'dual nationality' element, α \square .

Remark 2 The Theorem's statements benefit from brevity. But it is the Corollary's forms that most directly translate to a Linear-Algebraic formulation of the cycle of such results over all possible separations.

5 The first few nontrivial examples

Example 3

a) reads

$$3\mu_{\alpha} = \alpha + 2(\beta + \gamma) . \quad (19)$$

b) gives

$$M_A = \frac{B + C - A}{4} \quad (20)$$

I.e.

$$m_a^2 = \frac{b^2 + c^2 - a^2}{4} :$$

a standard form for Apollonius' Theorem recovered. For m_a the median corresponding to side a .

Example 4 [30] recovered.

a) reads

$$2 (\tau_1 + \tau_2) = \iota_1 - \alpha . \quad (21)$$

b) reads

$$\frac{2}{3} T_1 + \frac{3}{4} T_2 = \frac{1}{4} (S_1 - A) . \quad (22)$$

For which a rational form is

$$8 T_1 + 9 T_2 = 3 (S_1 - A) . \quad (23)$$

Or, in terms of the original Geometrical variables,

$$\frac{2}{3} t_1^2 + \frac{3}{4} t_2^2 = \frac{1}{4} (s_2 - a^2) . \quad (24)$$

With rational form

$$8 t_1^2 + 9 t_2^2 = 3 (s_2 - a^2) . \quad (25)$$

Example 5 For P_5 ,

a) reads

$$5 (\tau_1 + \tau_2 + \tau_3) = 2 \iota_1 - 3 \alpha . \quad (26)$$

b) reads

$$\frac{2}{3} T_1 + \frac{3}{4} T_2 + \frac{4}{5} T_3 = \frac{1}{10} (2 S_1 - 3 A) . \quad (27)$$

For which a rational form is

$$40 T_1 + 45 T_2 + 48 T_3 = 12 S_1 - 18 A . \quad (28)$$

Or, in terms of the original Geometrical variables,

$$\frac{2}{3} t_1^2 + \frac{3}{4} t_2^2 + \frac{4}{5} t_3^2 = \frac{1}{10} (2 s_2 - 3 a^2) . \quad (29)$$

With rational form

$$40 t_1^2 + 45 t_2^2 + 48 t_3^2 = 12 S_1 - 18 a^2 . \quad (30)$$

End Remark [31] go on to consider further eigenclustering that are not straight paths.

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